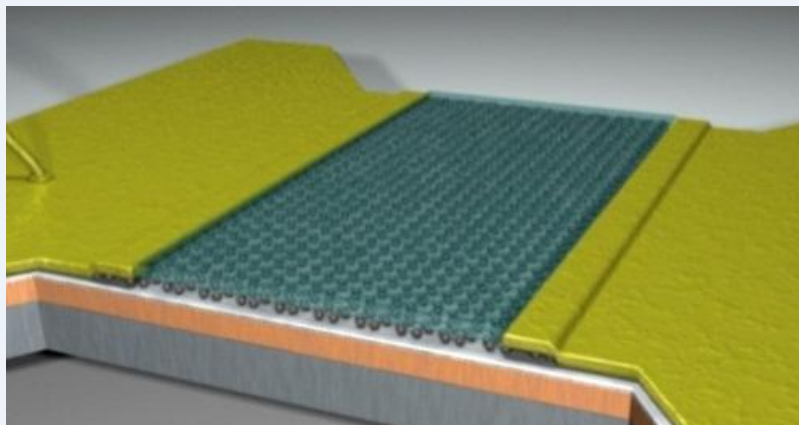


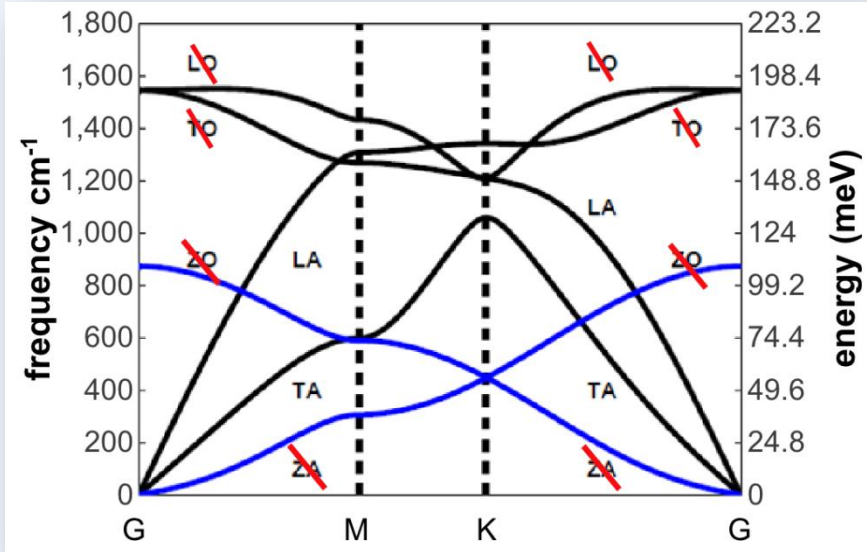
# Hot electron cooling by acoustic phonons in graphene



**Andreas Betz**, F. Vialla, D. Brunel, C. Voisin, G. Fève,  
J.-M. Berroir, B. Plaçais, E. Pallecchi (Ecole Normale Supérieure)  
M. Picher, A. Cavanna, A. Madouri  
(Laboratoire de Photonique et Nanostructures)

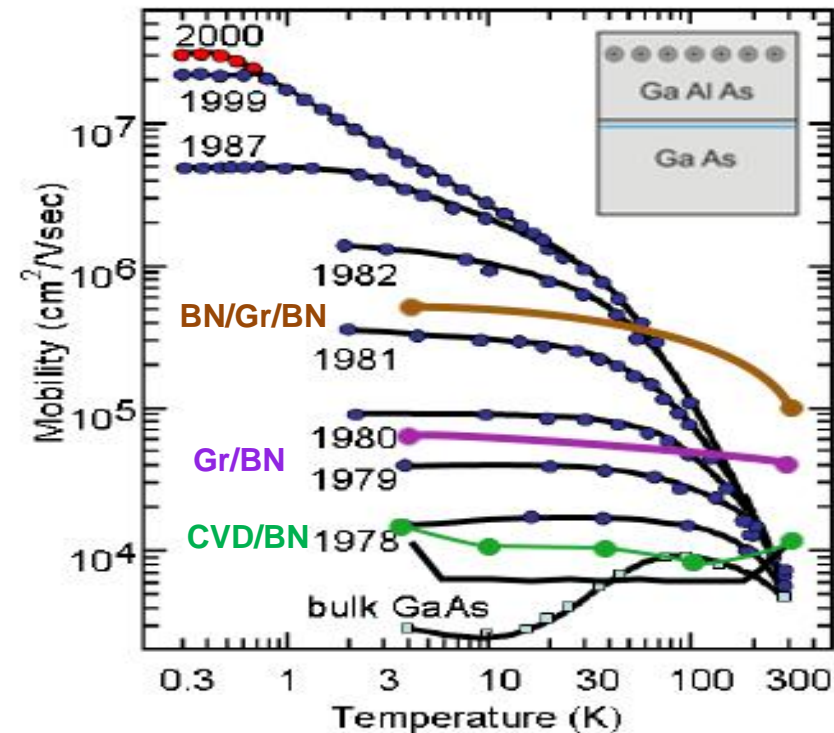
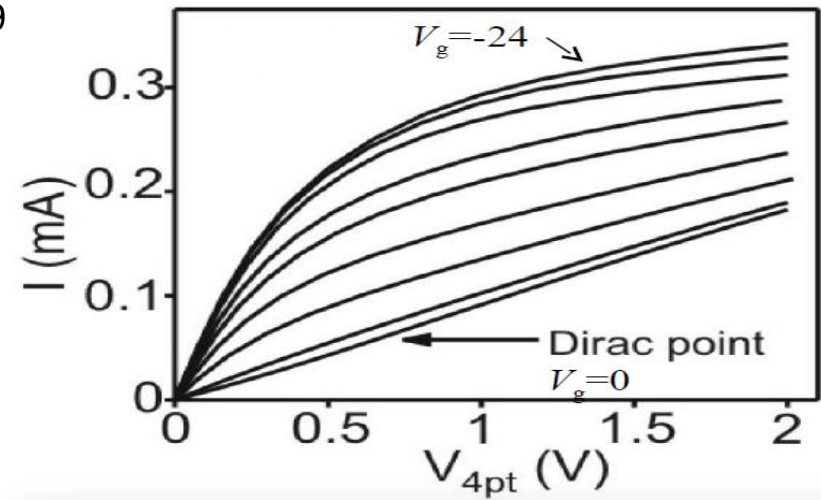
Laboratoire Pierre Aigrain – Ecole Normale Supérieure  
24 rue Lhomond, 75231 Paris Cedex 05 France

Barreiro, PRL 2009



optical  
 →  
 strong  
 interaction

acoustic  
 →  
 weak  
 interaction



■ large OP energy & strong interaction

■ small AP energy & weak interaction

Mayorov, Nano Lett., 2011, **11** (6)  
 Dean, Nature Nanotech **5**, 722–726 (2010)  
 Gannett, APL **98**, 242105(2011)

**What is the signature of 2D electron-acoustic phonon coupling?**

**How weak is acoustic phonon coupling in graphene?**

Theory :

Kubakaddi, PRB 2009

Viljas, Heikkila PRB 2010

Bistrizer, PRL 2009

Experiments :

Fay et al. PRB 2011 (noise thermometry)

Efetov et al. PRL 2010 (conductivity)

- 1) Electron-phonon cooling at 2D**
- 2) Fabrication & noise thermometry**
- 3) Results**

## Heat transfer to AP in the ideal metallic regime (high $n_s$ )

cooling power

(metals)

$$Q = \text{Volume} \times \Sigma \left( T_e^5 - T_{ph}^5 \right) \quad (3D)$$

$$Q = \text{Area} \times \Sigma \left( T_e^4 - T_{ph}^4 \right) \quad (2D)$$

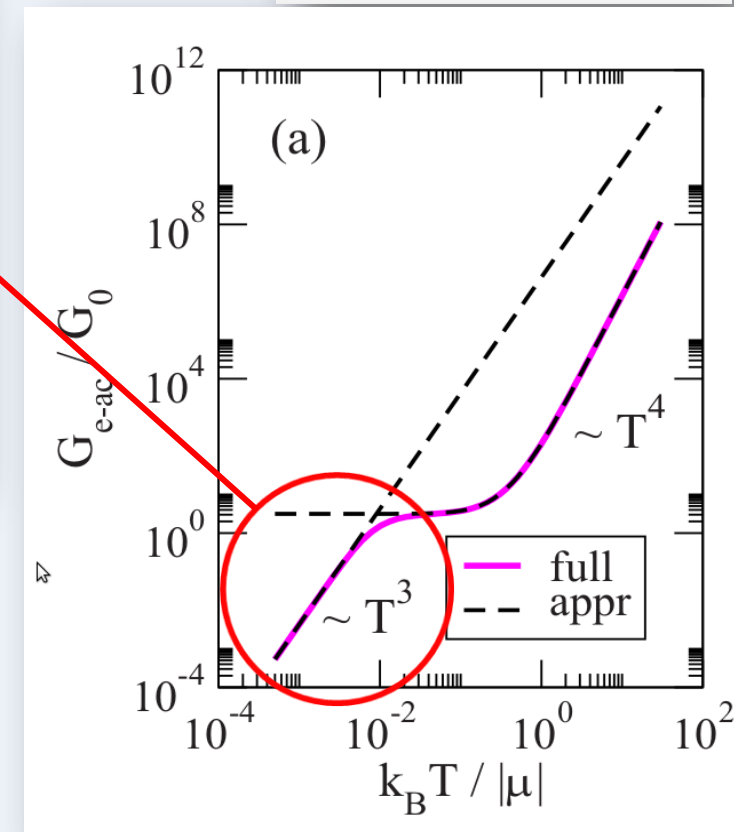
$$Q = \text{Length} \times \Sigma \left( T_e^3 - T_{ph}^3 \right) \quad (1D)$$

(nanotubes)

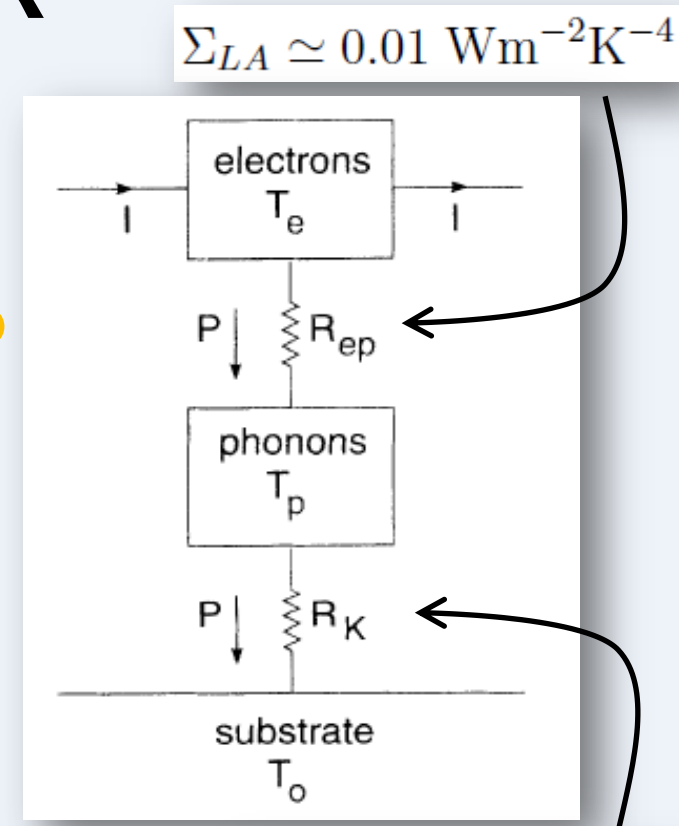
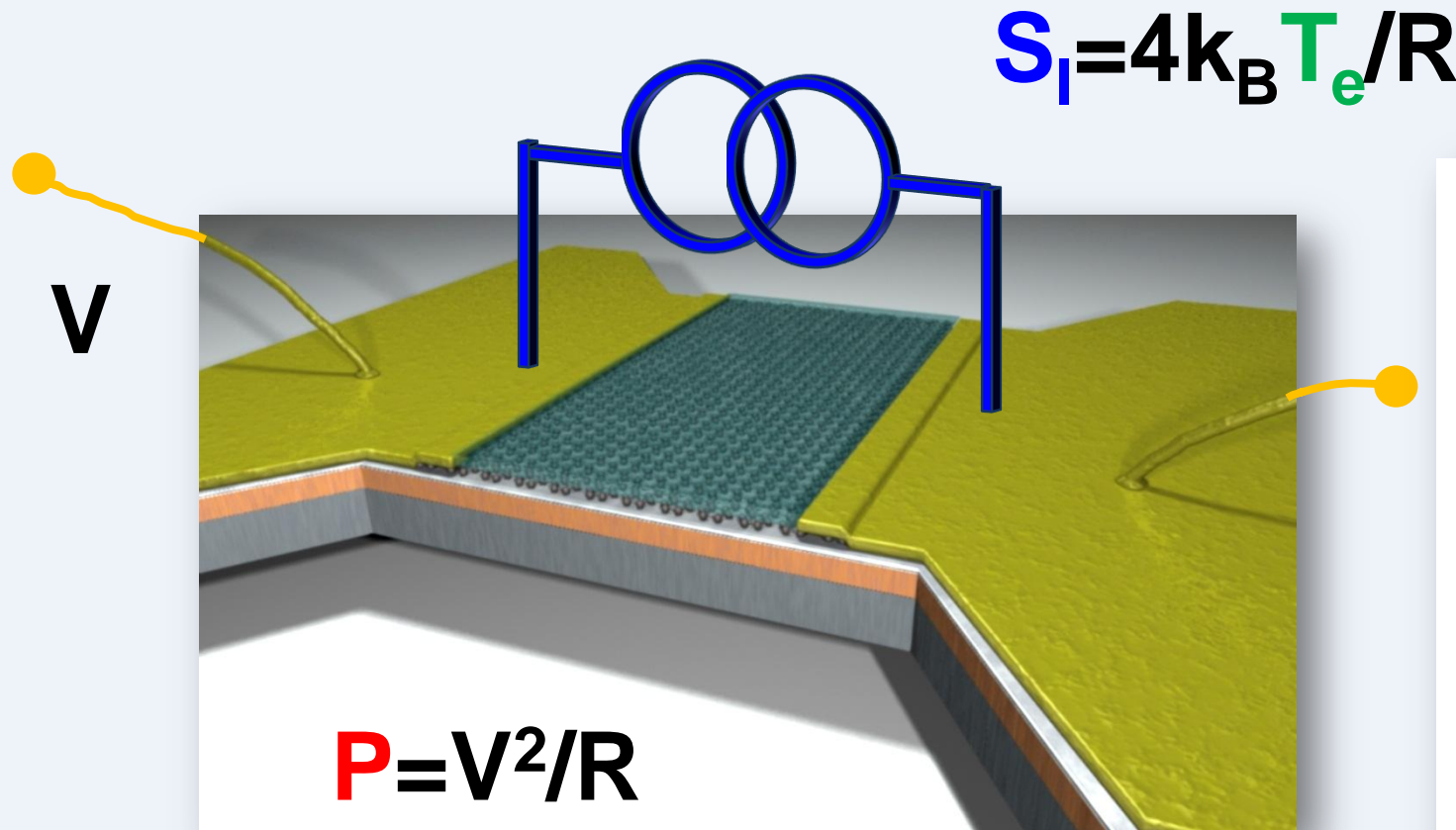
$$\Sigma_{LA} = \frac{\pi^2 D^2 k_B^4}{15 \rho \hbar^5 v_F^3 c^3} \times |E_F|$$

heat sink

$$G = 4 \Sigma T^3 \Delta T$$



Viljas, Heikkila PRB 2010



Direct measurement of AP coupling

Electron cooling power



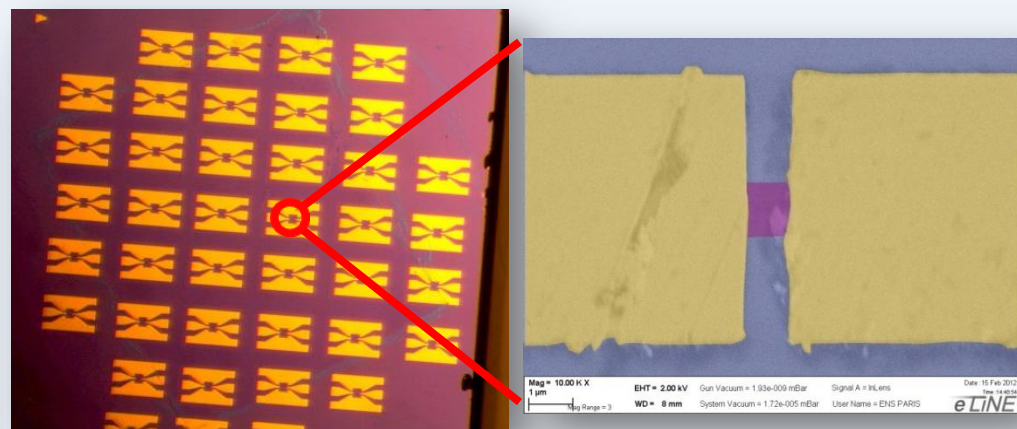
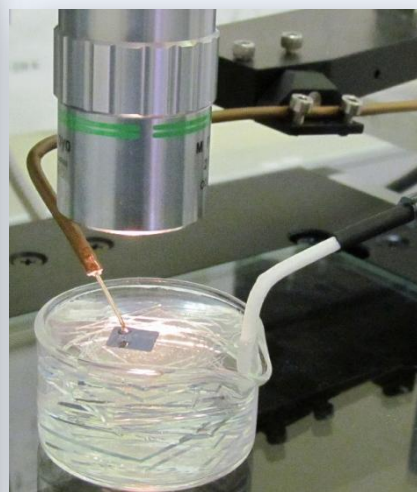
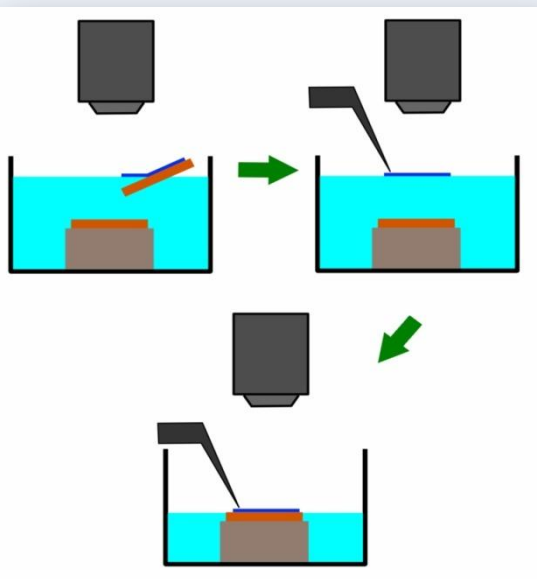
$$P = LW \Sigma (T_e^4 - T_p^4)$$

Balandin, Nature Mat. 2011

## Gr/BN

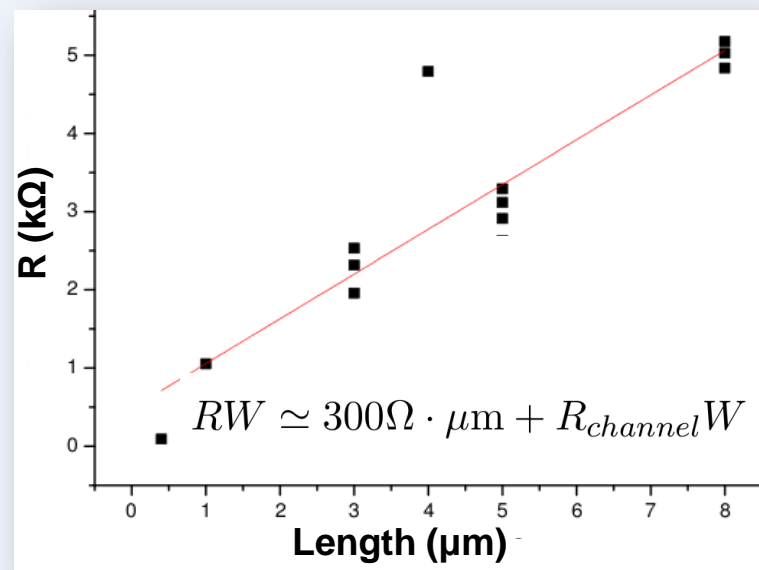
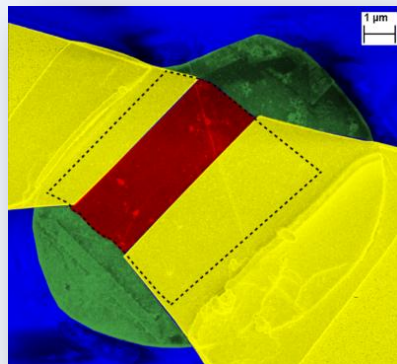
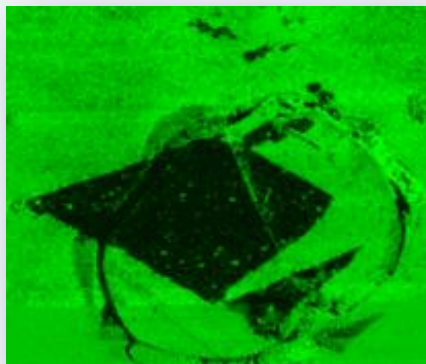
coplanar waveguide

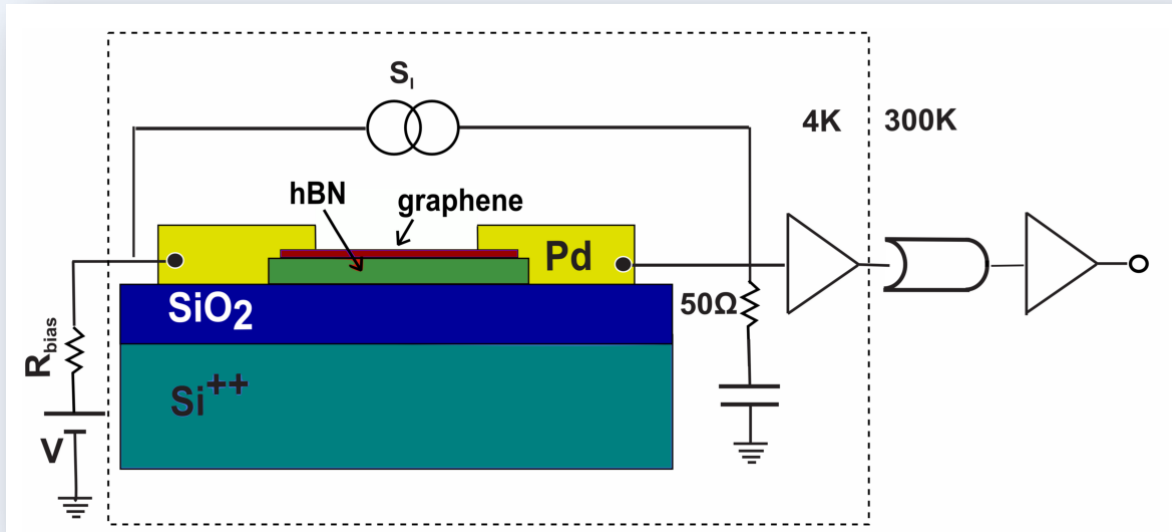
## CVD



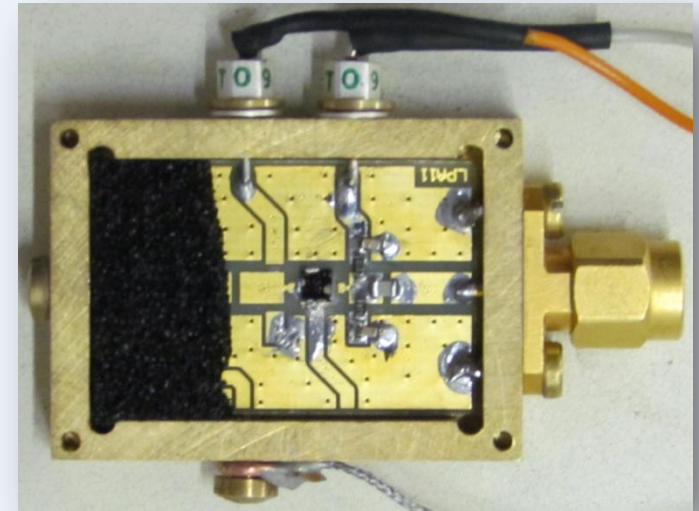
CVD grown on Cu

Wedging transfer: polymer CAB & water

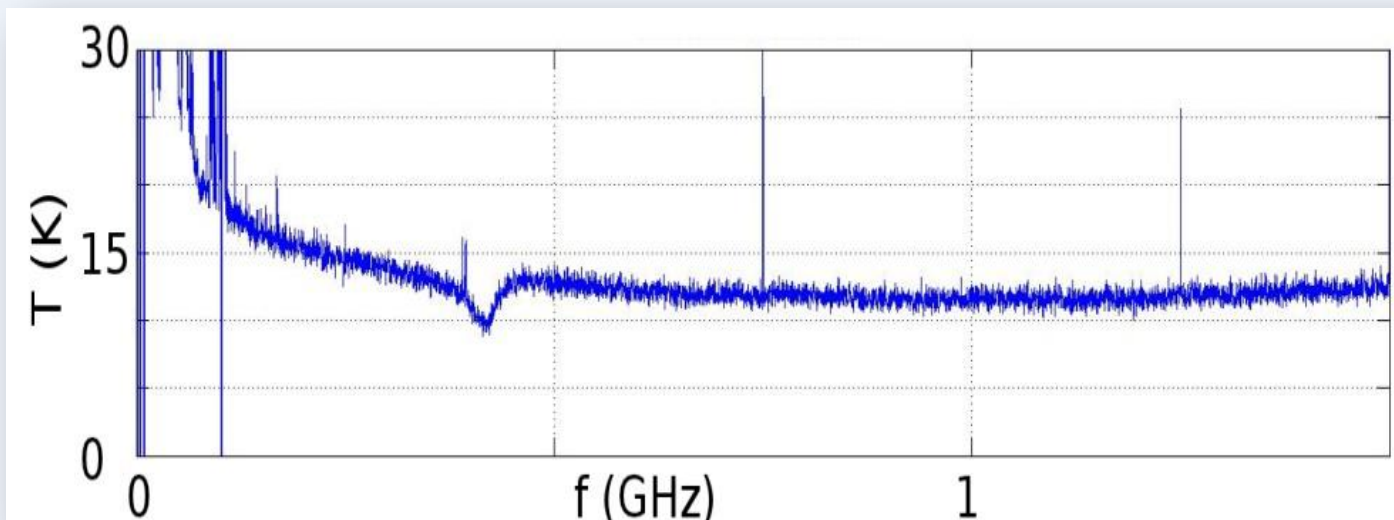




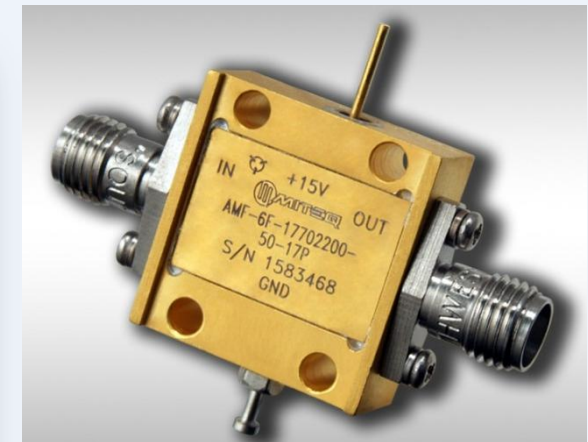
broadband GHz setup in liquid helium



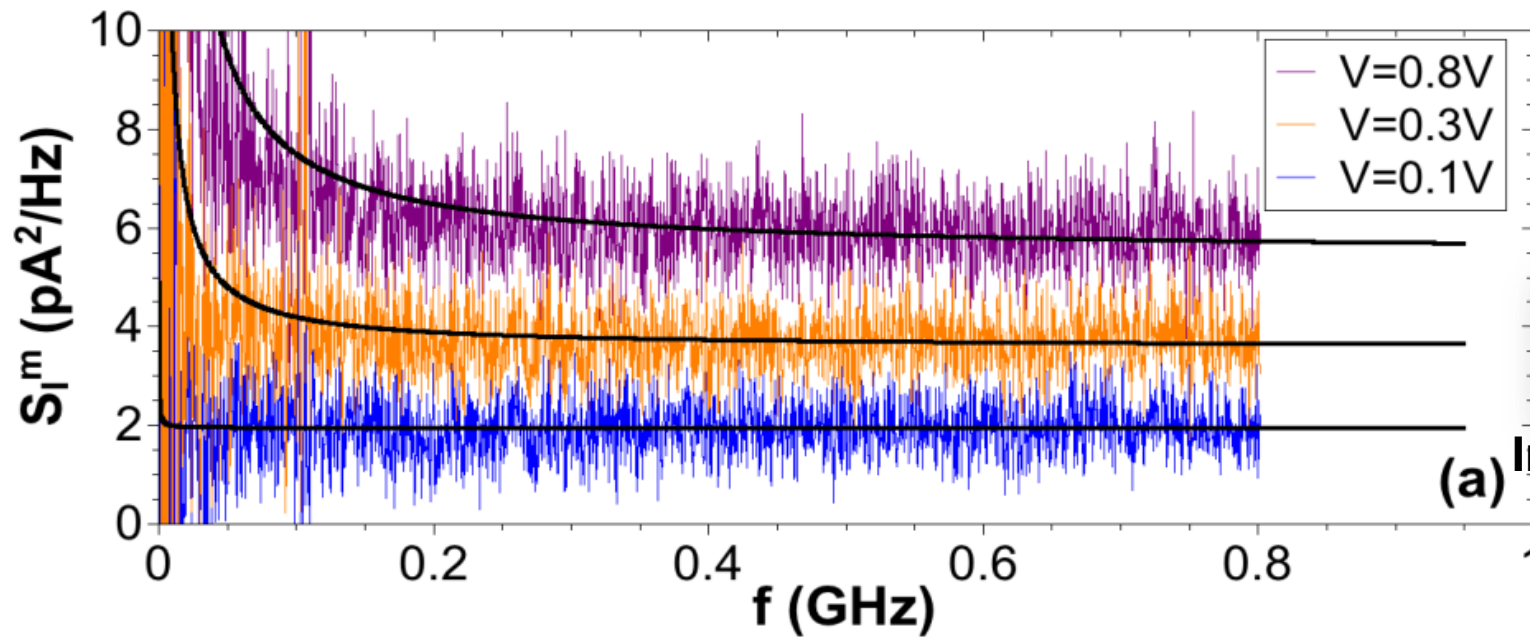
calibration: shot noise of an Al-AlOx-Al tunnel junction



Cryogenic LNA

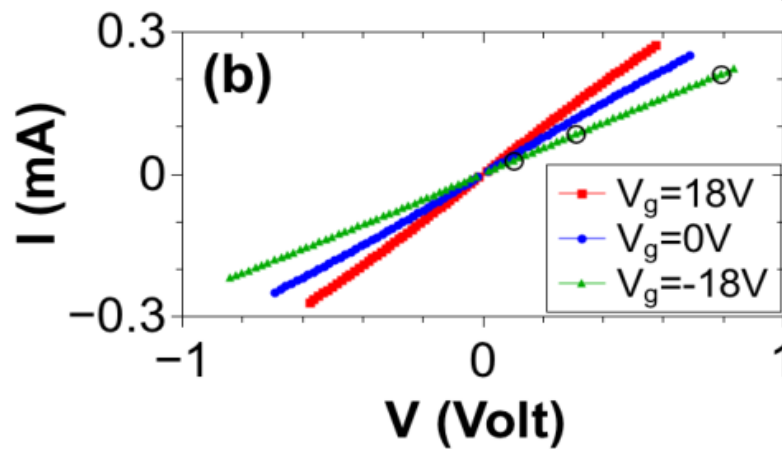




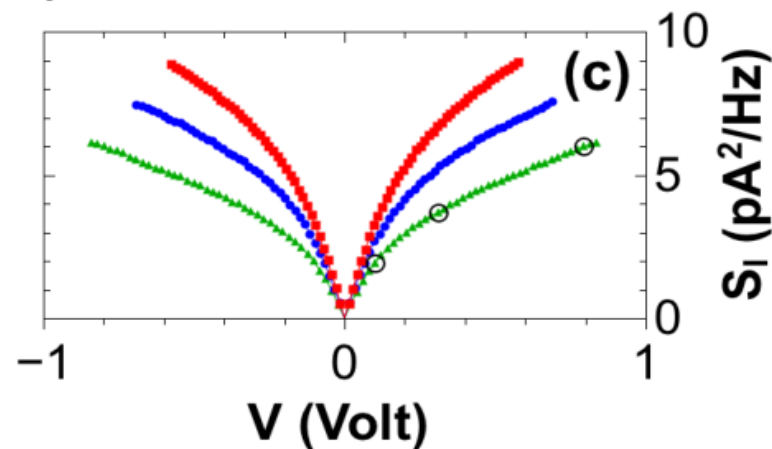


$$S_I^m \propto \frac{C \cdot I^2}{f} + S_I$$

Important for high bias  
& small samples

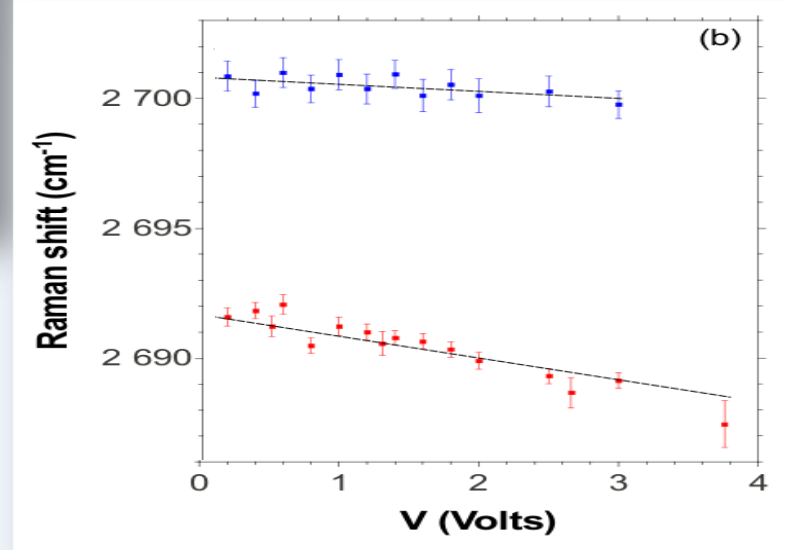
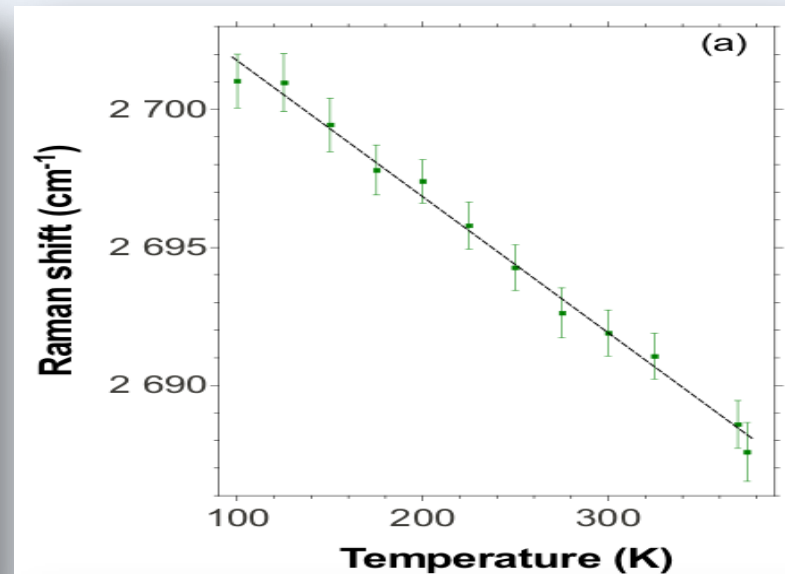
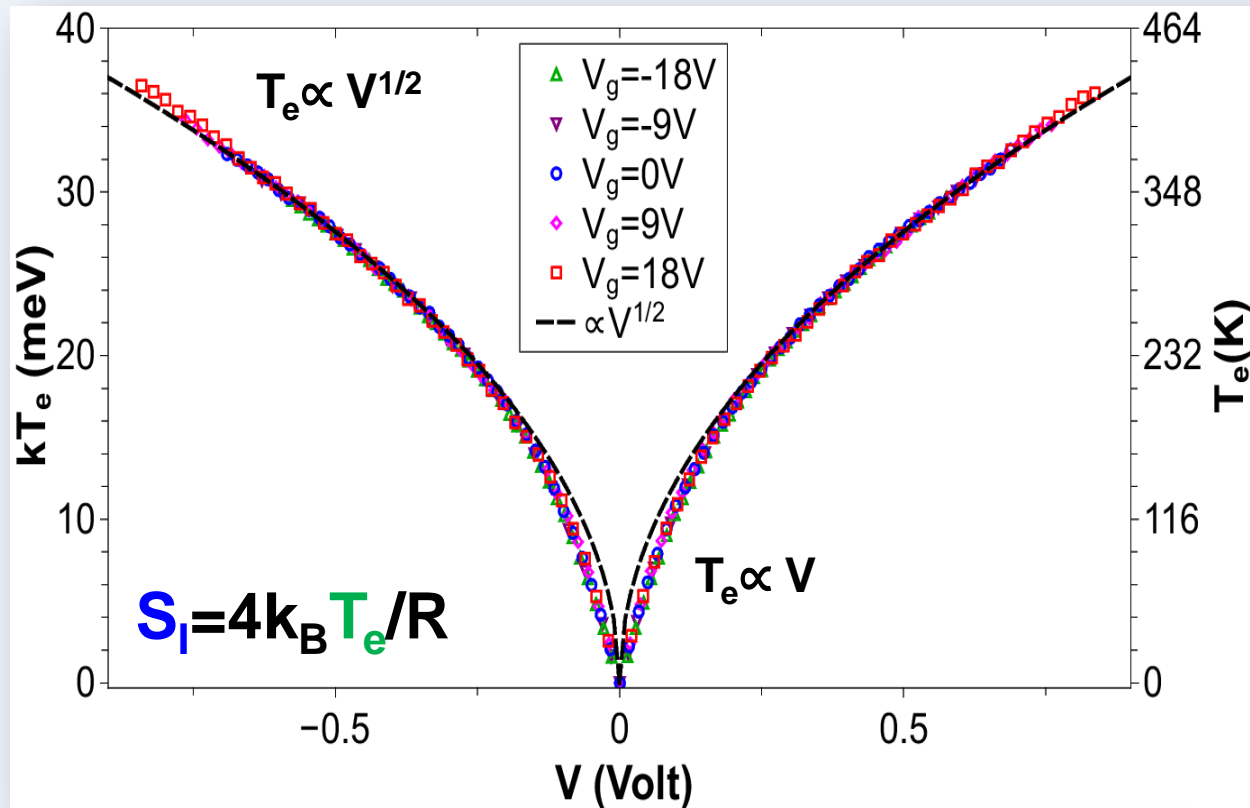


linear I-V : metallic regime



noise sub-linear

mobility = 350 - 3000 cm<sup>2</sup>/Vs



$$P = \frac{V^2}{R} \simeq LW \Sigma \left( T_e^4 - T_{ph}^4 \right)$$

→ Hot electrons (~400K/V)

→ evidence for  $T^4$  dependence

Cold phonons (~30K/V)

# Solving the heat equation

e<sup>-</sup> heat diffusion

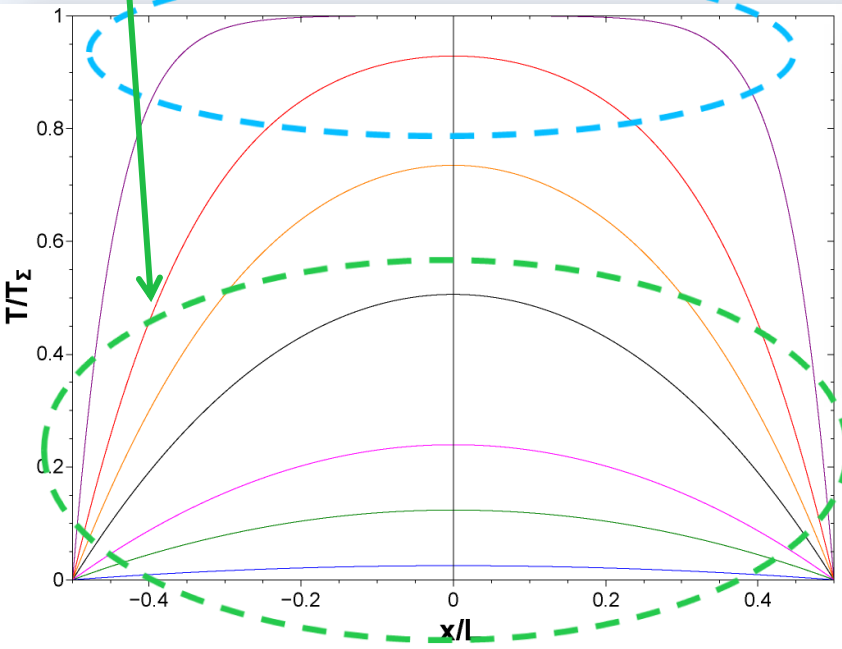
phonon cooling

$$\frac{\mathcal{L}}{2R} \frac{L^2 \partial^2 T^2(x)}{\partial x^2} = -\frac{V^2}{R} + LW\Sigma [T^4(x) - T_{ph}^4]$$

assumptions:

- uniform Joule heating
- cold contacts
- cold phonons

heat equation can be solved analytically:



Cold contacts ( $T(\pm L/2) = 0$ ) and cold phonons ( $T_{ph} = 0$ ) with  $U = (T/T_\Sigma)^2$

$$G(X) = \text{Elliptic} \left[ A \sin \sqrt{\frac{1}{2} + \frac{U_0 + 2X}{2\sqrt{3}\sqrt{4 - U_0^2}}}, \frac{2\sqrt{4 - U_0^2}}{\sqrt{3}U_0 + \sqrt{4 - U_0^2}} \right]$$

$$\frac{2x}{L} = \frac{[G(U_0) - G(U(x))]}{[G(U_0) - G(0)]} \quad ; \quad \frac{V}{V_\Sigma} = \frac{16 [G(U_0) - G(0)]^2}{U_0 + \sqrt{(4 - U_0^2)/3}}$$

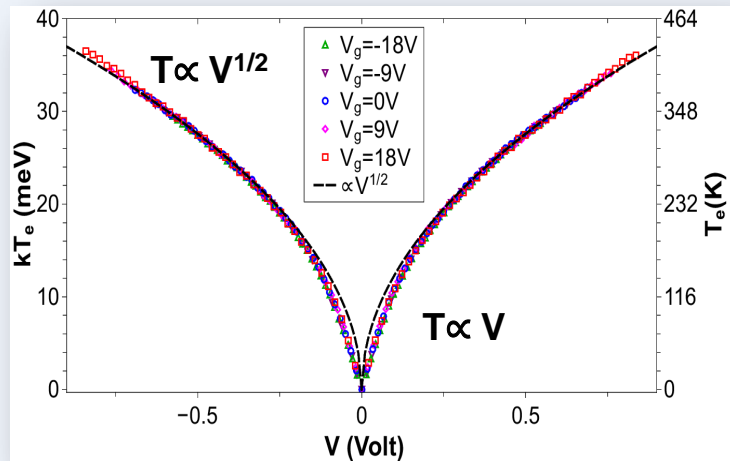
$$\text{T-scale : } T_\Sigma = \sqrt[4]{\frac{V^2}{RLW\Sigma}} = \sqrt[4]{\frac{P}{\Sigma}}$$

$$\text{V-scale : } V_\Sigma = \frac{\mathcal{L}}{\sqrt{4RLW\Sigma}}$$

→ global electron temperature  $\langle T_e \rangle$

$$S_I = \frac{4k_B}{R} \cdot \langle T_e \rangle$$

# The $P=\Sigma T^4$ dependence



$$\frac{\mathcal{L}}{2R} \frac{L^2 \partial^2 T^2(x)}{\partial x^2} = -\frac{V^2}{R} + LW\Sigma [T^4(x) - T_p^4]$$

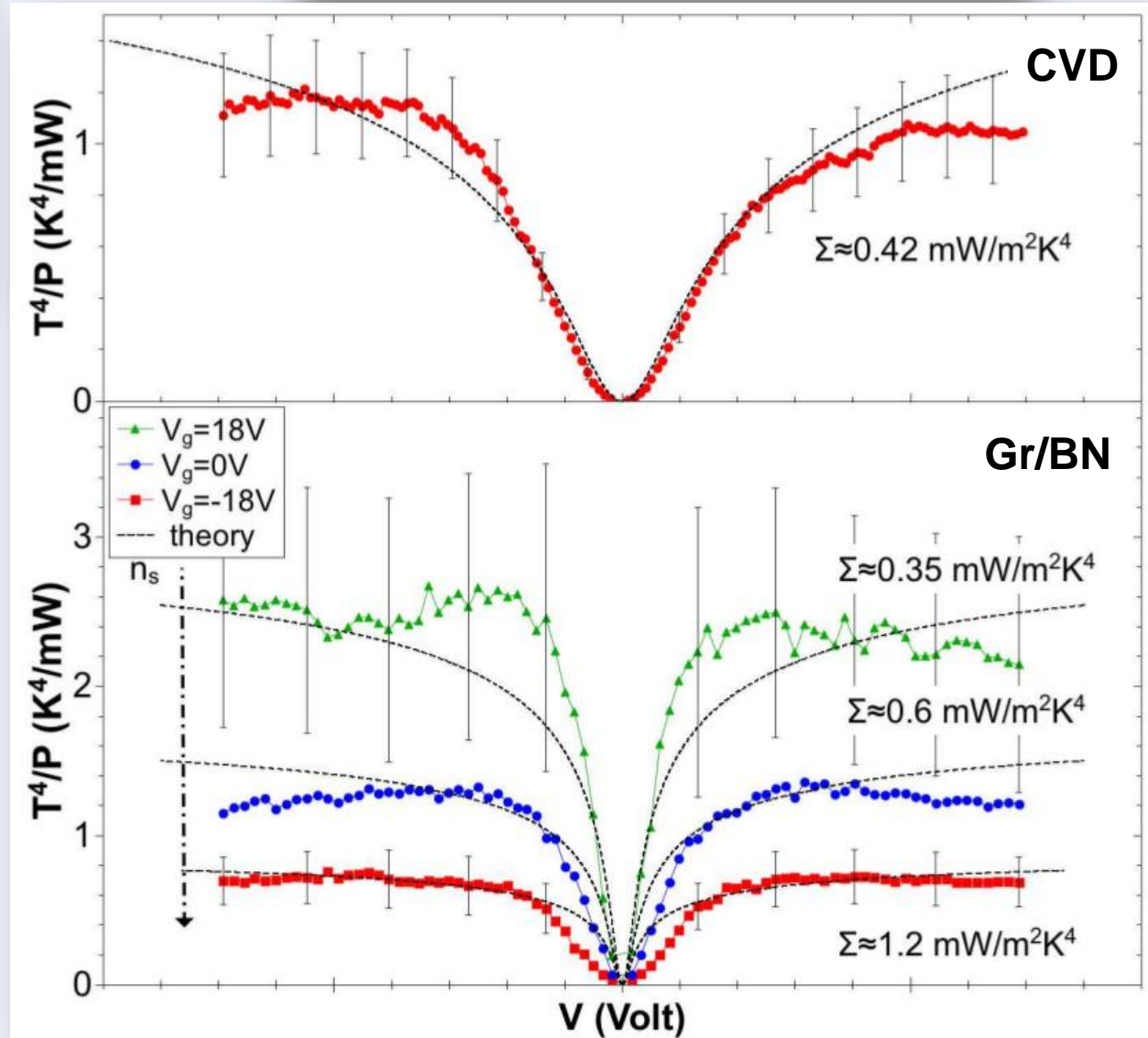
Use heat equation to fit  
ALL data (low & high V)

$$S_l = 4k_B T_e / R$$

$$P = I * V / (LW)$$

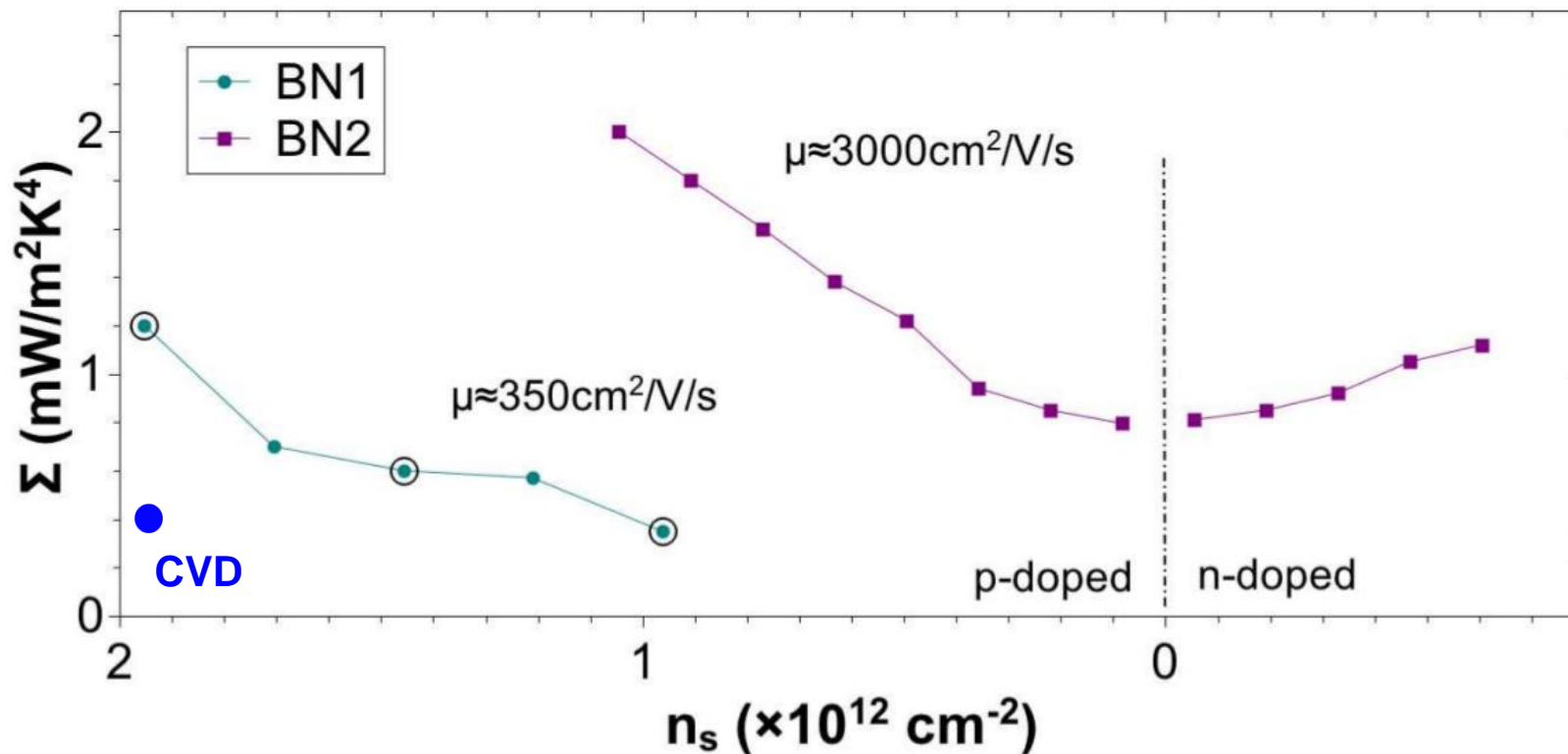
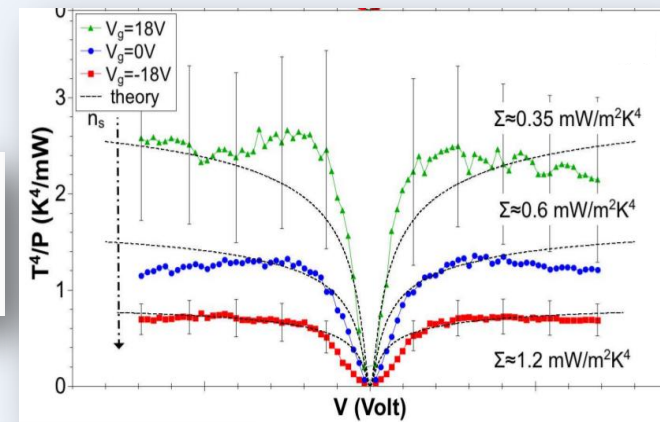
Normalise to power P

→ pronounce features:  
e- heat diffusion  
vs.  
phonon cooling



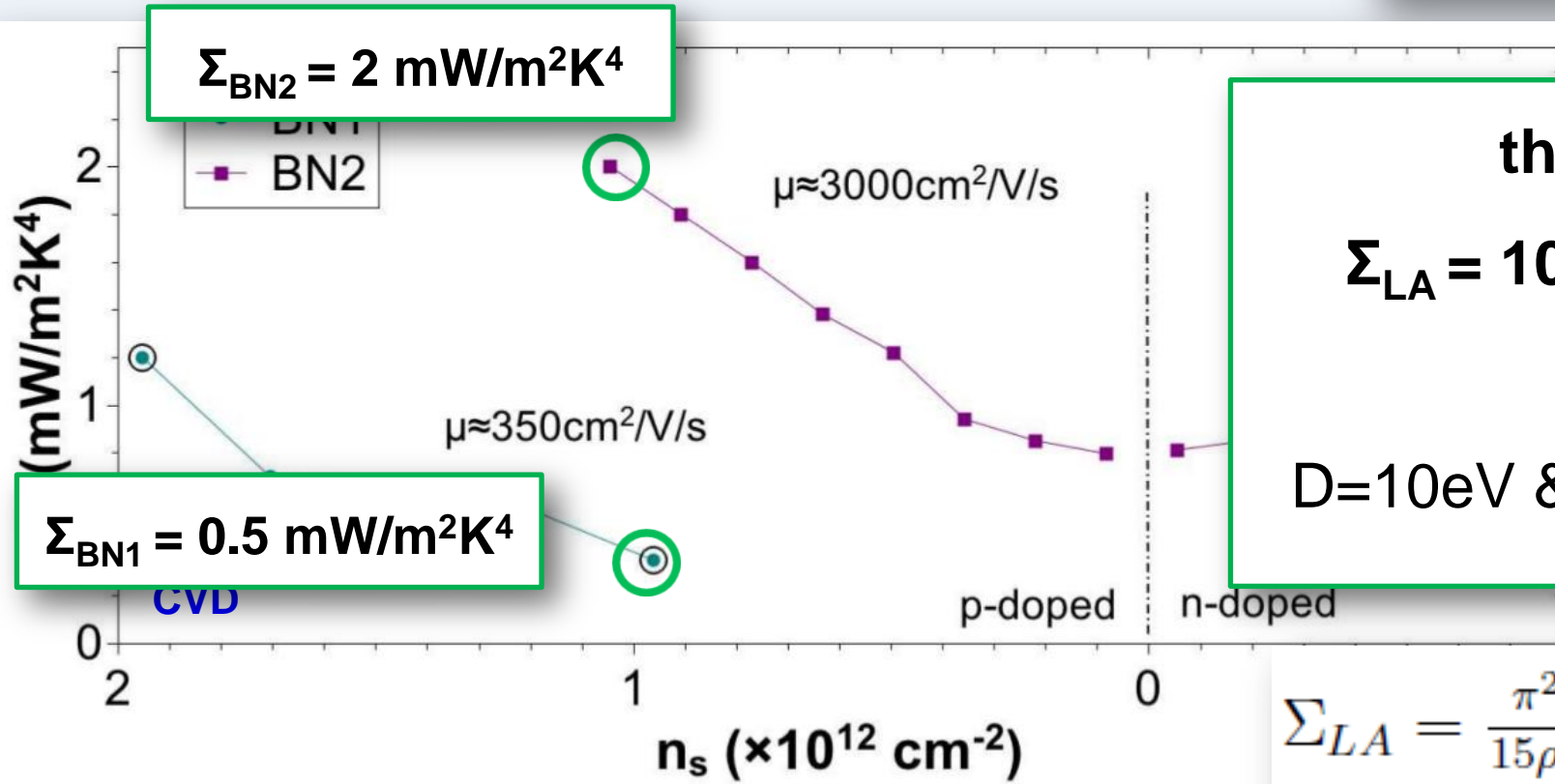
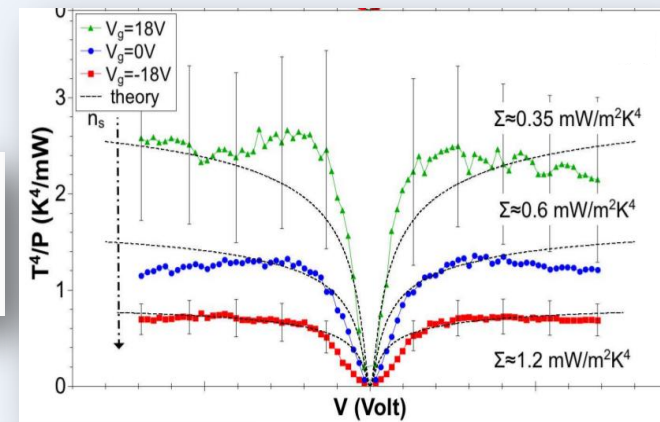
$\Sigma$  is the only free parameter in solution of

$$\frac{\mathcal{L}}{2R} \frac{L^2 \partial^2 T^2(x)}{\partial x^2} = -\frac{V^2}{R} + LW\Sigma [T^4(x) - T_{ph}^4]$$



$\Sigma$  is the only free parameter in solution of

$$\frac{\mathcal{L}}{2R} \frac{L^2 \partial^2 T^2(x)}{\partial x^2} = -\frac{V^2}{R} + LW\Sigma [T^4(x) - T_{ph}^4]$$



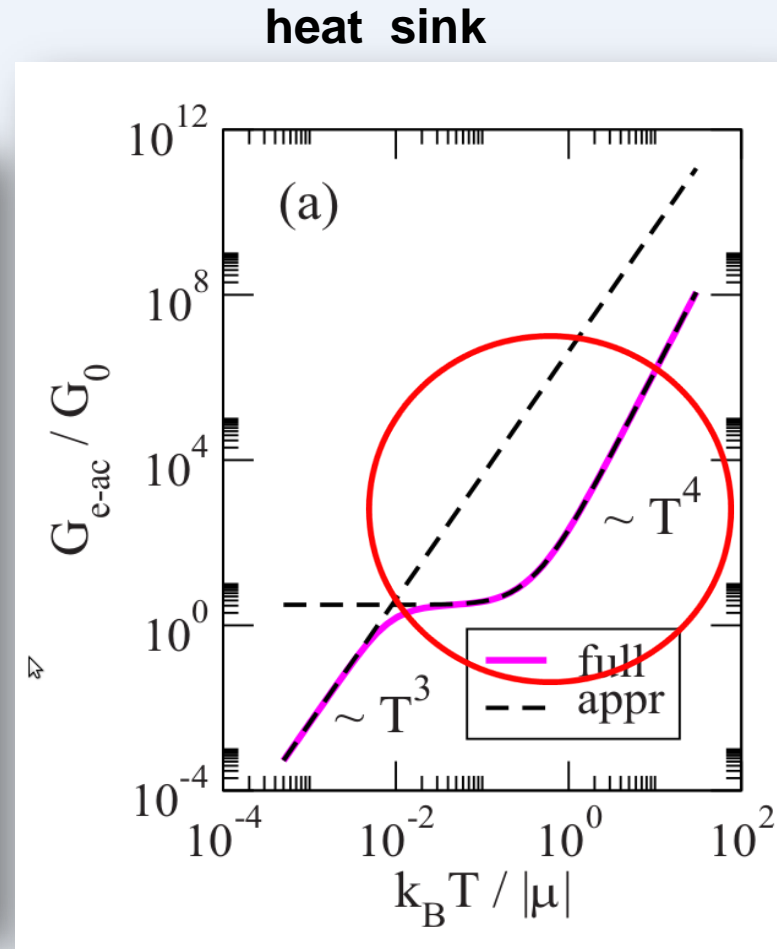
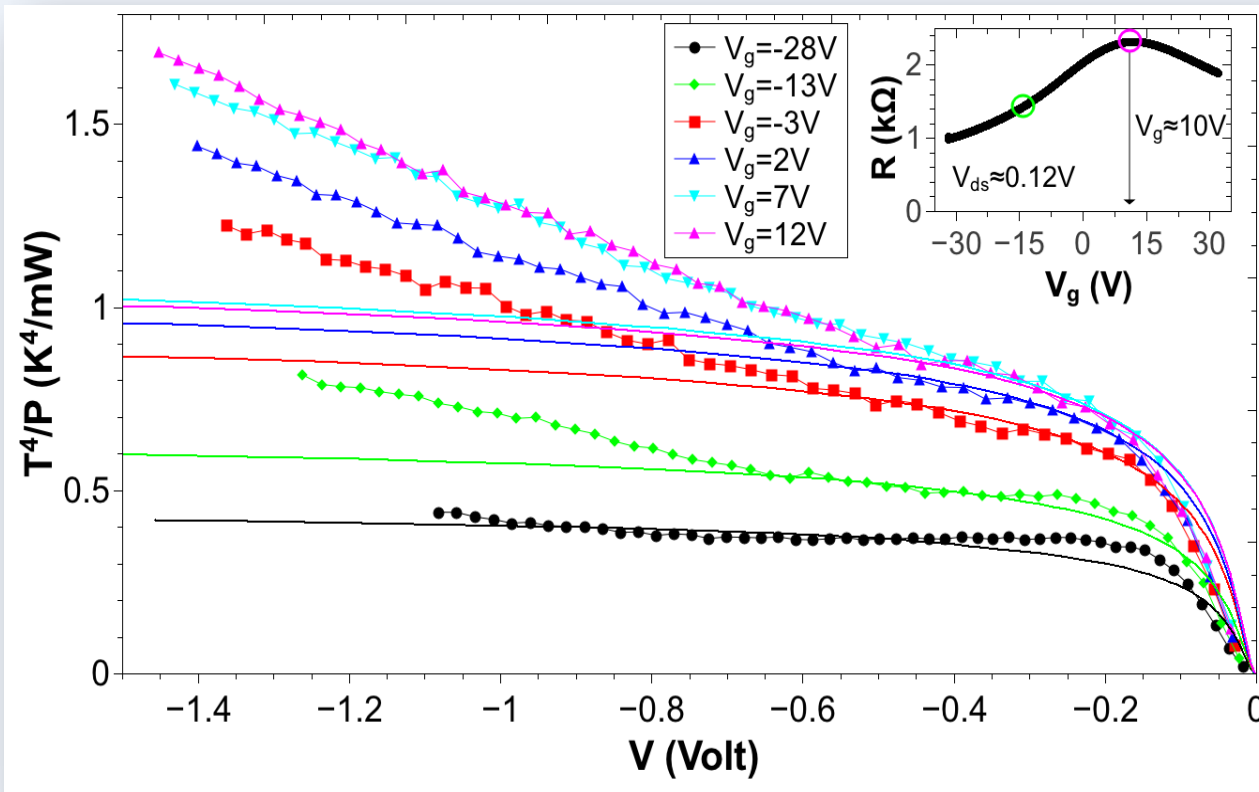
**theory:**

$\Sigma_{LA} = 10 \text{ mW/m}^2\text{K}^4$

for

$D=10\text{eV} \ \& \ n_s=10^{12}\text{cm}^{-2}$

$$\Sigma_{LA} = \frac{\pi^2 D^2 k_B^4}{15 \rho \hbar^5 v_F^3 c^3} \times |E_F|$$



*ongoing work*

- **GHz noise thermometry in graphene devices**
- **Qualitative signatures of acoustic phonon cooling**
- **Quantitative disagreement due to lattice disorder?**
- **Importance for graphene detectors**

Bolometric detection:

Gabor, Science 2011 (MIT : photo-current)

Vora, arxiv:1110.5623 (Stony Brook : 0.6 GHz)

Yan et al, arxiv:1111.1202 (Maryland : fast IR detector)

Fong, arxiv:1202.5737 (CALTECH : 1.2 GHz)

Freitag, arxiv:1202.5342 (IBM : photo-conductivity)



## Graphene team at LPA



Emiliano Pallecchi



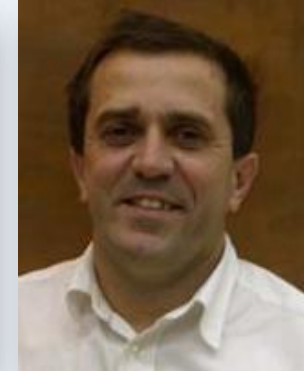
Andreas Betz



Jean-Marc Berroir



Gwendal Fève



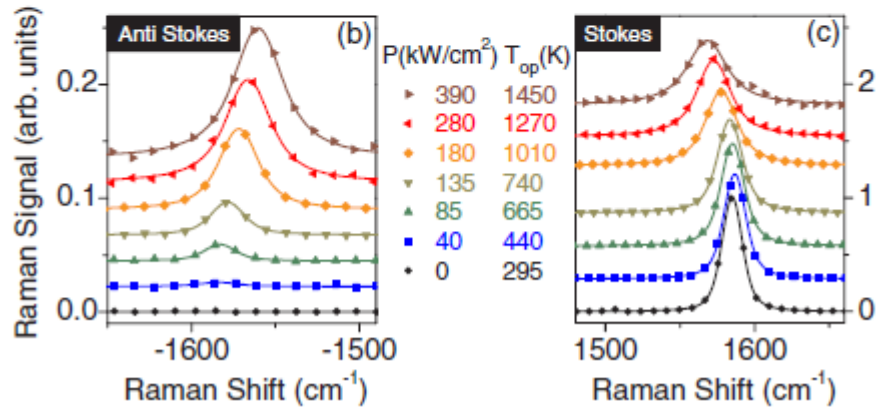
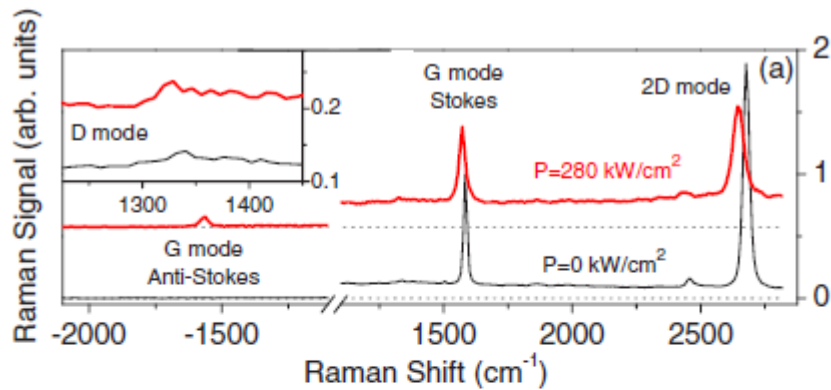
Bernard Plaçais

Laboratoire Pierre Aigrain – Ecole Normale Supérieure

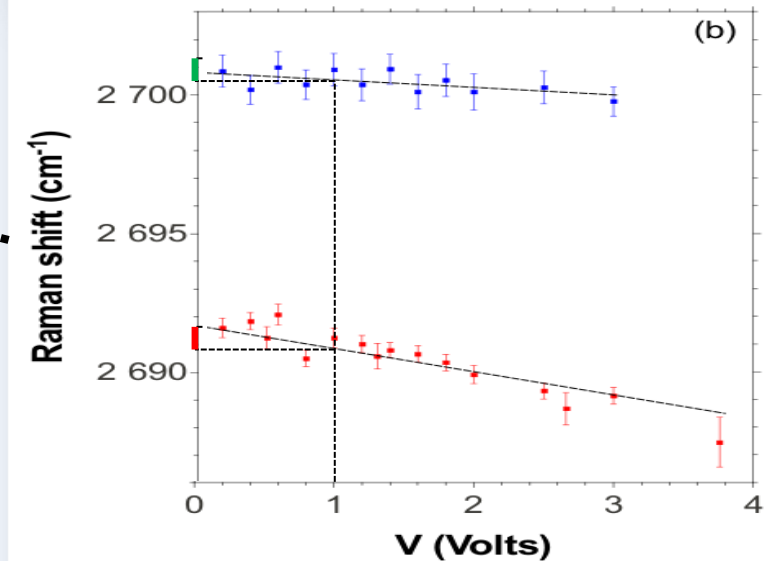
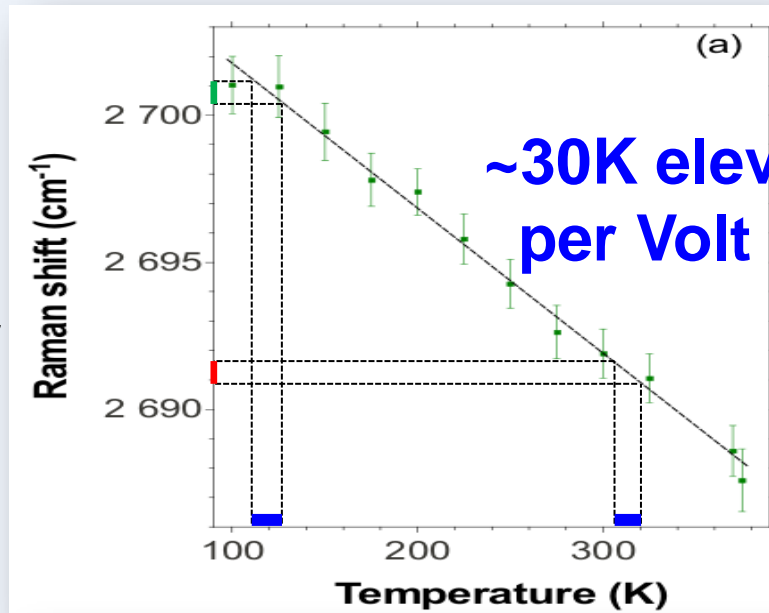
24 rue Lhomond, 75231 Paris Cedex 05 France

[www.lpa.ens.fr](http://www.lpa.ens.fr)

## OP/AP populations

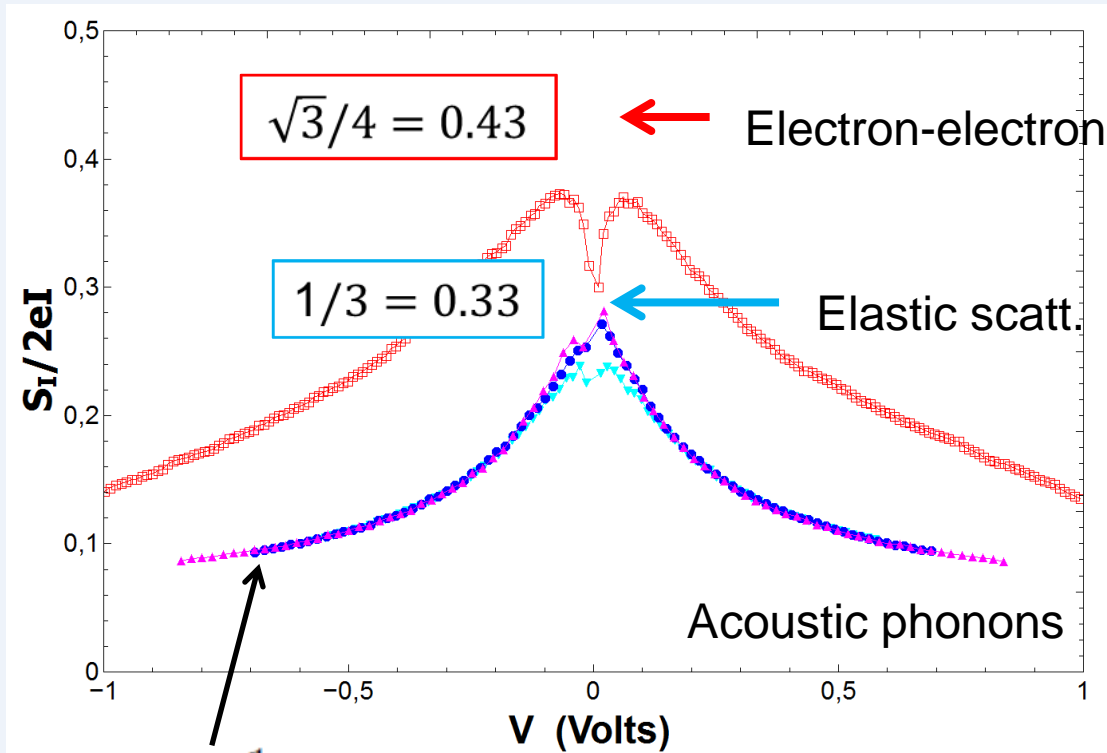


(Berciaud, Han, Mak, Brus, Kim, Heinz, PRL2010)



$$S_I/2eI \sim T_e/V$$

$$T_e/\sqrt{V}$$



$$F \propto \frac{1}{V}$$

